

Selected Exercises

Nuclear Forces PHY989

National Superconducting Cyclotron Laboratory and Department of Physics and Astronomy, Michigan State University, East Lansing, USA

Fall semester 2017

*

Exercise 1: Questions on the Overview of QCD

This collection of problems contain short exercises and discussion questions on QCD.

paragraphparagraph>paragraph>-0.5em

a) With respect to what scale(s) are the c, b, t quarks called heavy?

paragraphparagraph>paragraph>-0.5em

b) Have you heard about the s quark before? If yes, in what context?

paragraphparagraph>paragraph>-0.5em

c) A possible way to *see* quarks and gluons is in jets. What happens in these events?

paragraphparagraph>paragraph>-0.5em

d) Using the [Particle Data Group website](#), discuss which properties of the neutron and proton are similar and what are differences? What about for the three pions?

paragraphparagraph>paragraph>-0.5em

e) Which is more important in making a neutron more massive than a proton: the light quark mass difference or the electromagnetic contribution? Or do you think such considerations are too simplistic?

paragraphparagraph>paragraph>-0.5em

f) What is the evidence for *spontaneous* chiral symmetry breaking in

- the mass spectrum of pseudoscalar ($J^\pi = 0^-$) mesons;
- the mass spectrum of vector and axial vector ($J^\pi = 1^\mp$) mesons?

paragraphparagraph>paragraph>-0.5em

g) What is the evidence for *explicit* chiral symmetry breaking in the spectrum of pseudoscalar ($J^\pi = 0^-$) mesons?

paragraphparagraph>paragraph>-0.5em

h) If you and your friend each do a QCD calculation with the same diagrams but use α_s at different scales, will you get the same answer? If not, how could that happen?

paragraphparagraph>paragraph>-0.5em

i) Does the running coupling in QCD mean that the QCD Hamiltonian is not unique? Would you say that if you used α_s at two different scales that you were using two different Hamiltonians?

paragraphparagraph>paragraph>-0.5em

j) If the neutron lifetime is so short, why are there *any* stable nuclei?

paragraphparagraph>paragraph>-0.5em

k) One observes a marked resonance when a π^+ pion is scattering off a proton. Which baryon does this correspond to and at which energy of the π^+ does this occur (the proton is at rest)?

paragraphparagraph>paragraph>-0.5em

l) At sufficient energy in proton-proton collisions it is possible to create a pion, $p + p \rightarrow p + n + \pi^+$. At which energy in the center-of-mass frame does pion production start?

*

Exercise 2: Basic Scattering Theory

paragraphparagraph>paragraph>-0.5em

a) We typically use units in which $\hbar = c = 1$ and express quantities as powers of MeV or fm or both, using $\hbar c \approx 197.33$ MeVfm to convert between them. If we take for the nucleon mass $M_N = 939$ MeV/ c^2 , what is \hbar^2/M_N numerically in terms of MeV and fm?

Hint. Hint: This should be almost immediate if you insert the right factors of c .

paragraphparagraph>paragraph>-0.5em

b) For the scattering of equal mass (nonrelativistic) particles, if the laboratory energy E_{lab} is related to the magnitude of the relative momentum k_{rel} (i.e., the momentum each particle has in the center-of-mass frame) by $E_{\text{lab}} = Ck_{\text{rel}}^2$, what is C ? If the mass is $M_N = 939$ MeV, what is the value of C in MeVfm²?

paragraphparagraph>paragraph>-0.5em

c) We write the partial-wave momentum space Schroedinger equation (see Lecture notes) as

$$\frac{k^2}{2\mu} \langle klm|\psi\rangle + \frac{2}{\pi} \sum_{l'm'} \int_0^\infty dk' k'^2 \langle klm|V|k'l'm'\rangle \langle k'l'm'|\psi\rangle = E_k \langle klm|\psi\rangle ,$$

what are the units of $V_{ll'}(k, k') \equiv \langle klm|V|k'l'm\rangle$? In coordinate space the potential is local, $V(r)$, with units of MeV, and k is given in inverse fm. If you see a plot in a journal article of $V_{ll'}(k, k')$ with units of fm, how would you convert it to the units you just found?

Hint. Hint: use the results from the first exercise here.

paragraphparagraph>paragraph>-0.5em

d) In Figure 18 of the review by Scott Bogner *et al.*, [Prog. Nucl. Part. Phys.](#) **65**, 94 (2010) the momentum-space matrix elements of different chiral effective field theory potentials are given in units of fm. Consider the value at zero relative momenta. \tilde{C}_{1S_0} , see Eq. (2.5) and the [article by Epelbaum et al.](#) in GeV^{-2} . How do you convert to fm units? Do the values for the matrix elements then match?

paragraphparagraph>paragraph>-0.5em

e) What do *on-shell* and *off-shell* mean in the context of scattering?

paragraphparagraph>paragraph>-0.5em

f) Under what conditions is a partial-wave expansion of the potential useful?

paragraphparagraph>paragraph>-0.5em

g) Derive the standard result:

$$\frac{e^{i\delta_l(k)} \sin \delta_l(k)}{k} = \frac{1}{k \cot \delta_l(k) - ik}$$

paragraphparagraph>paragraph>-0.5em

h) Given a potential that is not identically zero as $r \rightarrow \infty$ (e.g., a Yukawa), how would you know in practice where the asymptotic (large r) region starts?

paragraphparagraph>paragraph>-0.5em

i) What is the physical interpretation of the relation between the (partial-wave) S -matrix and the scattering amplitude? (Note that $S_l(k) = 1 + 2ikf_l(k)$.)

*

Exercise 3: More on the Lippmann-Schwinger equation

paragraphparagraph>paragraph>-0.5em

a) Using the Schrödinger equation for the scattering of two particles with mass m ,

$$(H_0 + V)|\psi_E\rangle = E|\psi_E\rangle,$$

where H_0 is the free Hamiltonian, show that the Lippmann-Schwinger equation for the wave function,

$$|\psi_E^\pm\rangle = |\phi_k\rangle + \frac{1}{E - H_0 \pm i\epsilon} V |\psi_E^\pm\rangle,$$

is satisfied. Here $E = k^2/m$ and the plane wave state satisfies $H_0|\phi_k\rangle = E|\phi_k\rangle$. Why do you need the $\pm i\epsilon$?

paragraphparagraph>paragraph>-0.5em

b) We can define the T -matrix on-shell as the transition matrix that acting on the plane wave state yields the same result as the potential acting on the full scattering state. That is, $T^{(\pm)}(E = k^2/m)|\phi_k\rangle = V|\psi_E^\pm\rangle$. What does it mean that the T -matrix is *on-shell*? (This is a really quick question!)

paragraphparagraph>paragraph>-0.5em

c) Show that matrix elements of the T -matrix satisfy the Lippmann-Schwinger equation

$$\langle \mathbf{k}' | T^{(\pm)}(E) | \mathbf{k} \rangle = \langle \mathbf{k}' | V | \mathbf{k} \rangle + \int d^3p \frac{\langle \mathbf{k}' | V | \mathbf{p} \rangle \langle \mathbf{p} | T^{(\pm)}(E) | \mathbf{k} \rangle}{E - \frac{p^2}{m} \pm i\epsilon}.$$

What normalization is used for the momentum states? Are the matrix elements of the T -matrix on the right side on-shell?

paragraphparagraph>paragraph>-0.5em

d) Write the Lippmann-Schwinger equation for the wave function in coordinate space for a local potential $V = V(\mathbf{r})$. To this end, show first that the free Green's function

$$G^\pm(\mathbf{r}', \mathbf{r}; E = k^2/m) = \langle \mathbf{r}' | \frac{1}{E - H_0 \pm i\epsilon} | \mathbf{r} \rangle,$$

is given by

$$G^\pm(\mathbf{r}', \mathbf{r}; E = k^2/m) = -\frac{m}{4\pi} \frac{e^{\pm ik|\mathbf{r}-\mathbf{r}'|}}{|\mathbf{r}-\mathbf{r}'|}.$$

paragraphparagraph>paragraph>-0.5em

e) Show that when the T -matrix is evaluated on-shell, it is proportional to the scattering amplitude, $T^+(E = k^2/m) = -\frac{1}{4\pi^2 m} f(k, \theta)$, by analyzing the asymptotic form of the Lippmann-Schwinger equation and comparing to

$$\langle \mathbf{r} | \psi_E^+ \rangle \xrightarrow{r \rightarrow \infty} (2\pi)^{-3/2} \left(e^{i\mathbf{k} \cdot \mathbf{r}} + f(k, \theta) \frac{e^{ikr}}{r} \right).$$

paragraphparagraph>paragraph>-0.5em

f) Start from the momentum-space partial wave expansion of the potential,

$$\langle \mathbf{k}' | V | \mathbf{k} \rangle = \frac{2}{\pi} \sum_{l,m} V_l(k', k) Y_{lm}^*(\Omega_{k'}) Y_{lm}(\Omega_k),$$

and a similar expansion of the T -matrix to derive the partial wave version of the Lippmann-Schwinger equation (with the correct factor for the integral):

$$T_l(k', k; E) = V_l(k', k) + \frac{2}{\pi} \int_0^\infty dp p^2 \frac{V_l(k', q) T_l(q, k; E)}{E - p^2/m + i\epsilon}.$$

paragraphparagraph>paragraph>-0.5em

g) Scattering phase shifts for a square well potential. Calculate the S-wave scattering phase shifts for an attractive square-well potential $V(r) = -V_0\theta(R-r)$ and show that

$$\delta_0(E) = \arctan \left[\sqrt{\frac{E}{E + V_0}} \tan(R\sqrt{2\mu(E + V_0)}) \right] - R\sqrt{2\mu E}.$$

paragraphparagraph>paragraph>-0.5em

h) Let's consider the analytic structure of the corresponding partial-wave S matrix, which is given by

$$S_0(k) = e^{-2ikR} \frac{k_0 \cot k_0 R + ik}{k_0 \cot k_0 R - ik},$$

where $E = k^2/2\mu$ and $k_0^2 = k^2 + 2\mu V_0$. Show that $S_l(k) = e^{2i\delta_l(k)}$ for $l = 0$ is satisfied. Treat $S_0(k)$ as a function of the complex variable k and find its singularities.

Hint. Hint: write $e^{2i\delta} = e^{i\delta}/e^{-i\delta}$.

paragraphparagraph>paragraph>-0.5em

i) Bound states are associated with poles on the imaginary axis in the upper half plane. Show that the condition for such a pole here gives the same eigenvalue condition (a transcendental equation) that you would get from a conventional solution to the square well by matching logarithmic derivatives.

Hint. Hint: Define $k = i\kappa$ with $\kappa > 0$ when analyzing such a pole.